



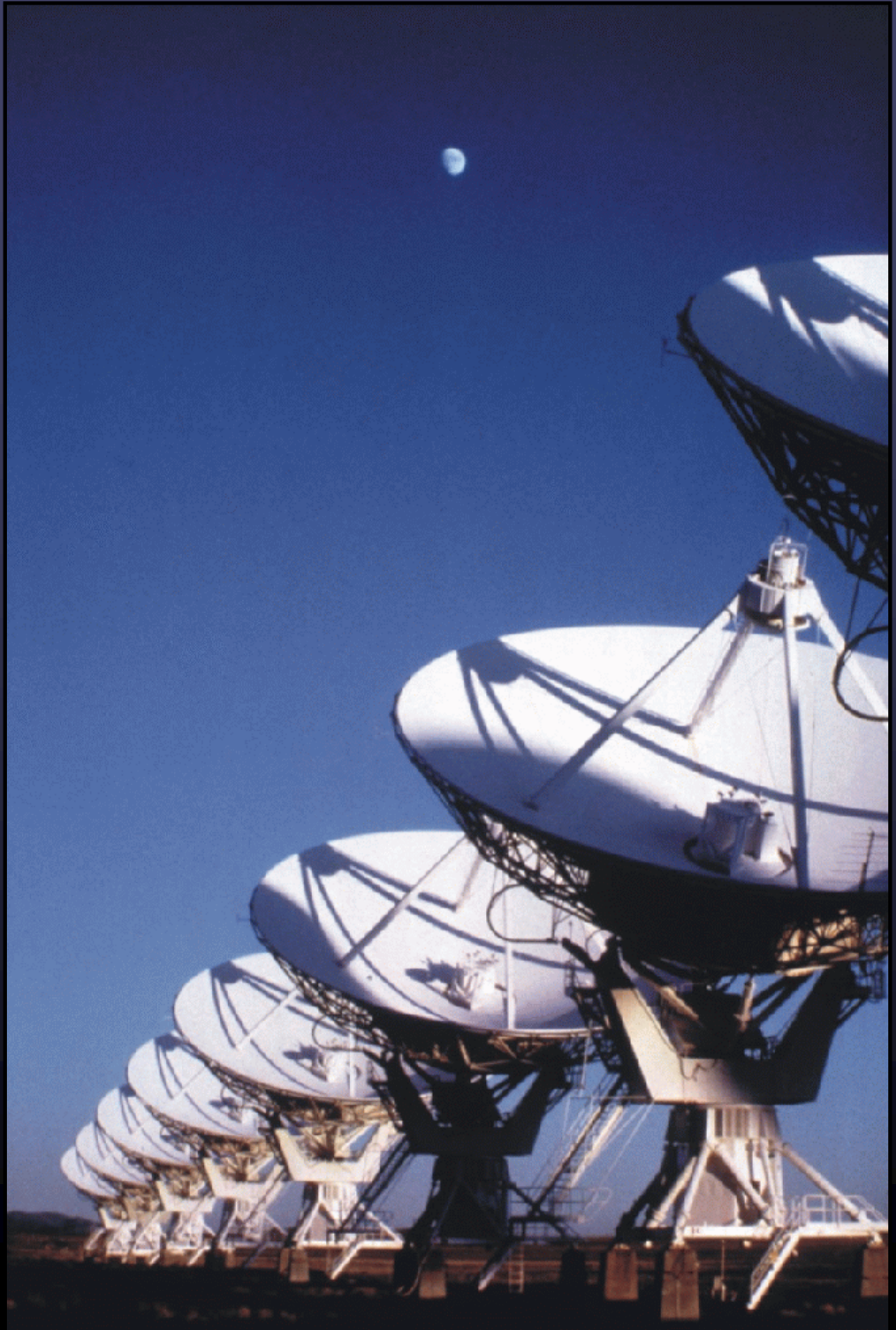
# Radio Astronomy Bremsstrahlung Radiation

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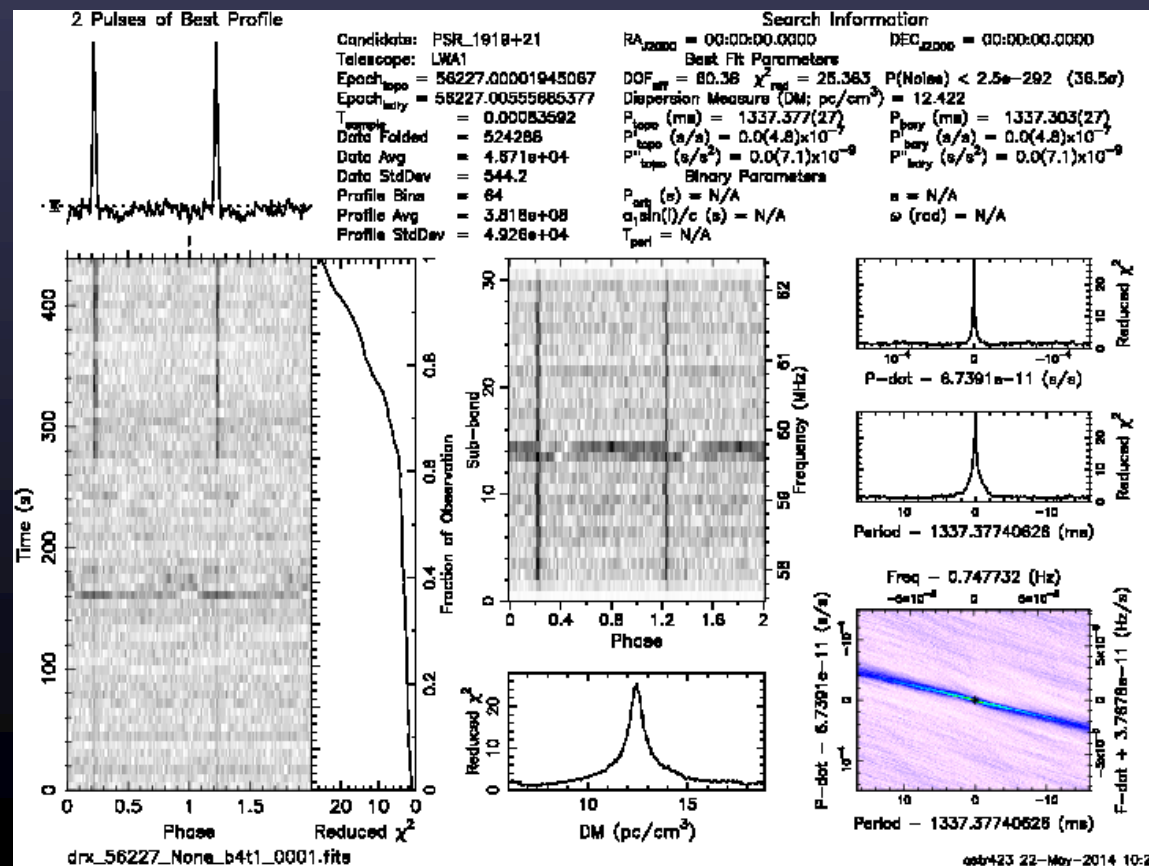
*Astronomy 423 at UNM*

*Radio Astronomy*



# Announcements

- Show pulsar observing on OPS page
- For HW#9 problem 4 you need to analyze the results of your LWA1 observations and submit the PDF file created by prepfold. See example:



## Announcements

- Where's my data?
- Connect to Hercules with x2go, then
  - ssh -X lwaucf2 (password is the same)
  - cd /data/local/astr423/axxx (xxx is group/session ID)
  - Data is in /data/network/recent\_data/astr423

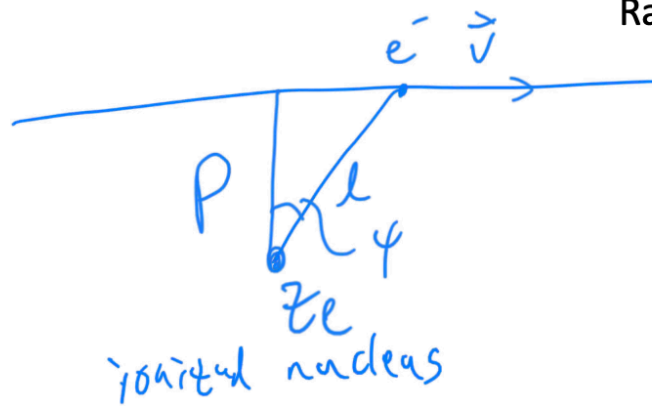
```
astr423@lwaucf2:/data/local/astr423$ ls /data/network/recent_data/astr423
059274_000047893 059278_000020808 059282_000059084 059316_000138250
059274_000240409 059278_000023487 059282_000067478 COMGT_0422.tgz
059276_000001376 059280_000040138 059316_000133529 COMGT_0423.tgz
059276_000001675 059280_000045300 059316_000134997 COMGT_0424.tgz
```

- Except for session ID 421 which observes soon



# Bremsstrahlung (free-free) radiation

Radio Astronomy Notes



assume motion of  $e^-$  is in a straight line (deflection is small)

$l = \frac{p}{\cos \phi}$       accel of  $e^-$  is given by Coulomb's law

$$F = ma = m\dot{v} = -\frac{Ze^2}{l^2}$$

$$\dot{v} = |\dot{v}| \cos \phi$$

$$\dot{v} = \frac{-Ze^2}{m_e p^2} \cos^3 \phi$$

Dipole       $w = \frac{\frac{4}{3} Z^2 e^6}{c^3 m_e p^4} \int_0^\phi \cos^6 \phi(t) dt$

would prefer as integral of  $\phi$

define  $dt = \frac{1}{2} l^2 d\psi$

$$\frac{dF}{dt} = \frac{1}{2} l^2 \frac{d\psi}{dt} = \text{const}$$

now solve for  $\frac{dF}{dt}$

$$\text{at } t=0 \quad \frac{dF}{dt} = \frac{1}{2} \rho x \frac{v}{\rho} = \frac{1}{2} \rho v$$

$$dt = \frac{l^2}{v \rho} d\psi = \frac{\rho}{v \cos^2 \psi} d\psi$$

$$W = \frac{4}{3} \frac{z^2 e^6}{c^3 m_e^2 \rho^4} \frac{\rho}{v} \int_0^{\frac{\pi}{2}} \cos^4 \psi d\psi \Rightarrow \frac{4}{16}$$

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$$W = \frac{\pi}{4} \frac{z^2 e^4}{c^3 m_e^2 p^3} \frac{1}{v}$$

total energy radiated by a single  $e^-$

What does spectrum look like?

$$A(\omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} E(t) e^{i\omega t} dt \quad \text{by FT}$$

$$C(\omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \dot{v}(t) \cos(\omega t) dt$$

$$= \frac{-ze^2}{m_e p^2} \frac{1}{2\pi} \int_0^{\infty} \cos(\omega t) \cos^3 \psi(t) dt$$





$$\omega_g = 0.463 \frac{v}{p}$$

radiation from a cloud of  $e^-$ s

maxwellian velocity distribution

$$f(v) = \frac{4v^2}{\sqrt{\pi}} \left( \frac{m}{2kT} \right)^{3/2} e^{-\frac{mv^2}{2kT}}$$

$$P_\nu(p, \nu) = \frac{16}{3} \frac{\pi^2 e^6}{c^3 m^2} \frac{1}{p^2 \nu^2} \quad \text{for } \nu < \nu_g \quad \text{and } 0 \text{ for } \nu > \nu_g$$

Ne in a volume  $^{-3}$  with collisions between p and p+dp

emissivity  $\int 4\pi p dp v N_e f(v) dv$

$\int 4\pi E_\nu dv = P_\nu(v, p) dN(v, p) dv$

$\int 4\pi E_\nu = \int_{p_1}^{p_2} \int_0^\infty \frac{8}{3} \frac{z^2 e^6}{c^3 m^2} \frac{1}{p^2 v^2} N_i N_e f(v) 2\pi p dp dv$

$E_\nu = \frac{8}{3} \frac{z^2 e^6}{c^3} \frac{N_i N_e}{m^2} \sqrt{\frac{2m}{\pi kT}} \ln\left(\frac{p_2}{p_1}\right)$

↑  
Gault factor

or

$E_\nu = 5.44 \times 10^{-39} g_{ff} \frac{z^2 n_e n_i}{T^{1/2}} e^{-h\nu/kT} \text{ erg cm}^{-2} \text{ s}^{-1} \text{ sr}^{-1} \text{ Hz}^{-1}$



recall

$$\frac{\tau_\nu}{K_\nu} = B_\nu(T) \quad \text{Planck Function}$$

for R.T.  $B_\nu(T) = \frac{2kT}{\lambda^2}$

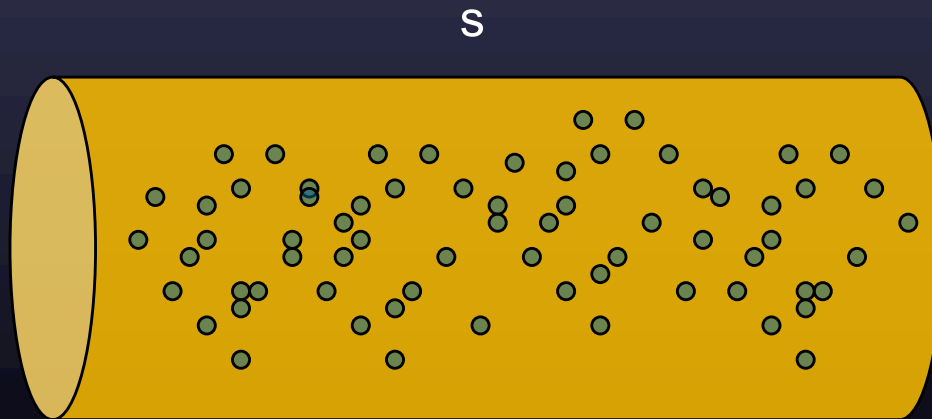
$$K_\nu = \frac{\tau_\nu}{B_\nu(T)} = \frac{4\pi^2 e^6}{3c} \frac{n_i n_e}{\nu^2} \frac{1}{\sqrt{2\pi(mkT)^3}} \ln\left(\frac{P_2}{P_1}\right)$$

assume  $n_i = n_e$  and uniform  $T$

$$\tau_\nu = - \int_0^S K_\nu ds$$

$$\tau_\nu = 3.014 \times 10^{-2} \left(\frac{T}{K}\right)^{-3/2} \left(\frac{\nu}{\text{GHz}}\right)^{-2} \left(\frac{EM}{\mu\text{C cm}^{-6}}\right) g_{ff}$$

Where  $EM = \int_0^S n_e^2 ds$        $n_e$  in  $\text{cm}^{-3}$   
 $S$  in  $\mu\text{C}$



Emission Measure (EM) is a column density

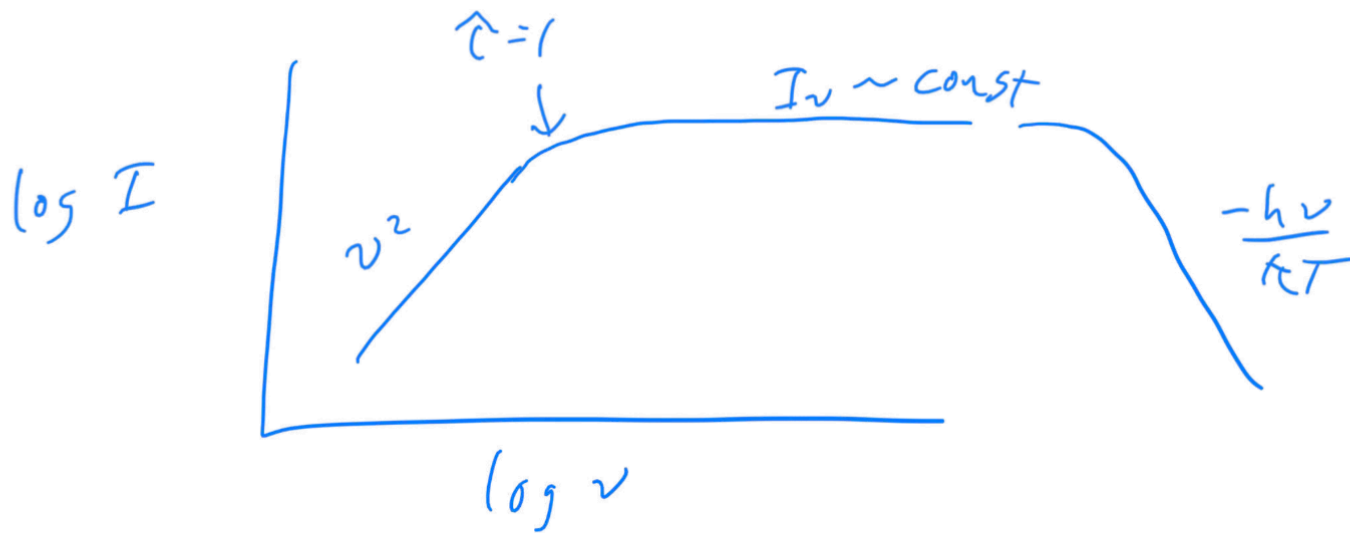
$$\tau_\nu = 3.014 \times 10^{-2} \left(\frac{T}{K}\right)^{-3/2} \left(\frac{\nu}{\text{GHz}}\right)^{-2} \left(\frac{EM}{\text{pc cm}^{-6}}\right) g_{ff}$$

where  $EM = \int_0^S n_e^2 ds$        $n_e$  in  $\text{cm}^{-3}$   
 $S$  in  $\text{pc}$

$$W = 1.435 \times 10^{-27} \tau^2 g_{ff} T^{1/2} n_e^2 \text{ erg cm}^{-3} \text{ s}^{-1}$$

correcting for  $g_{ff}$  (slow function of  $T$  and  $\nu$ )

$$\tau_\nu = 8.235 \times 10^{-2} \left(\frac{T_e}{K}\right)^{-1.35} \left(\frac{\nu}{\text{GHz}}\right)^{-2.1} \left(\frac{EM}{\text{pc cm}^{-6}}\right)$$



At high temperatures velocity distribution no longer Maxwellian and  $e^-$ s are relativistic in the tail.

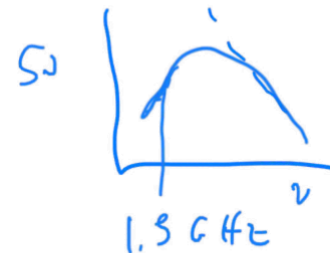
Example: Consider an AGN with absorption in a disk around the central engine

$T \sim 8000 \text{ K}$  path  $50 \text{ pc}$

$S_{\text{obs}} = 10 \text{ mJy}$  at  $1.5 \text{ GHz}$  what is  $n_e$ ?  
 $S_{\text{exp}} = 250 \text{ mJy}$  " " "

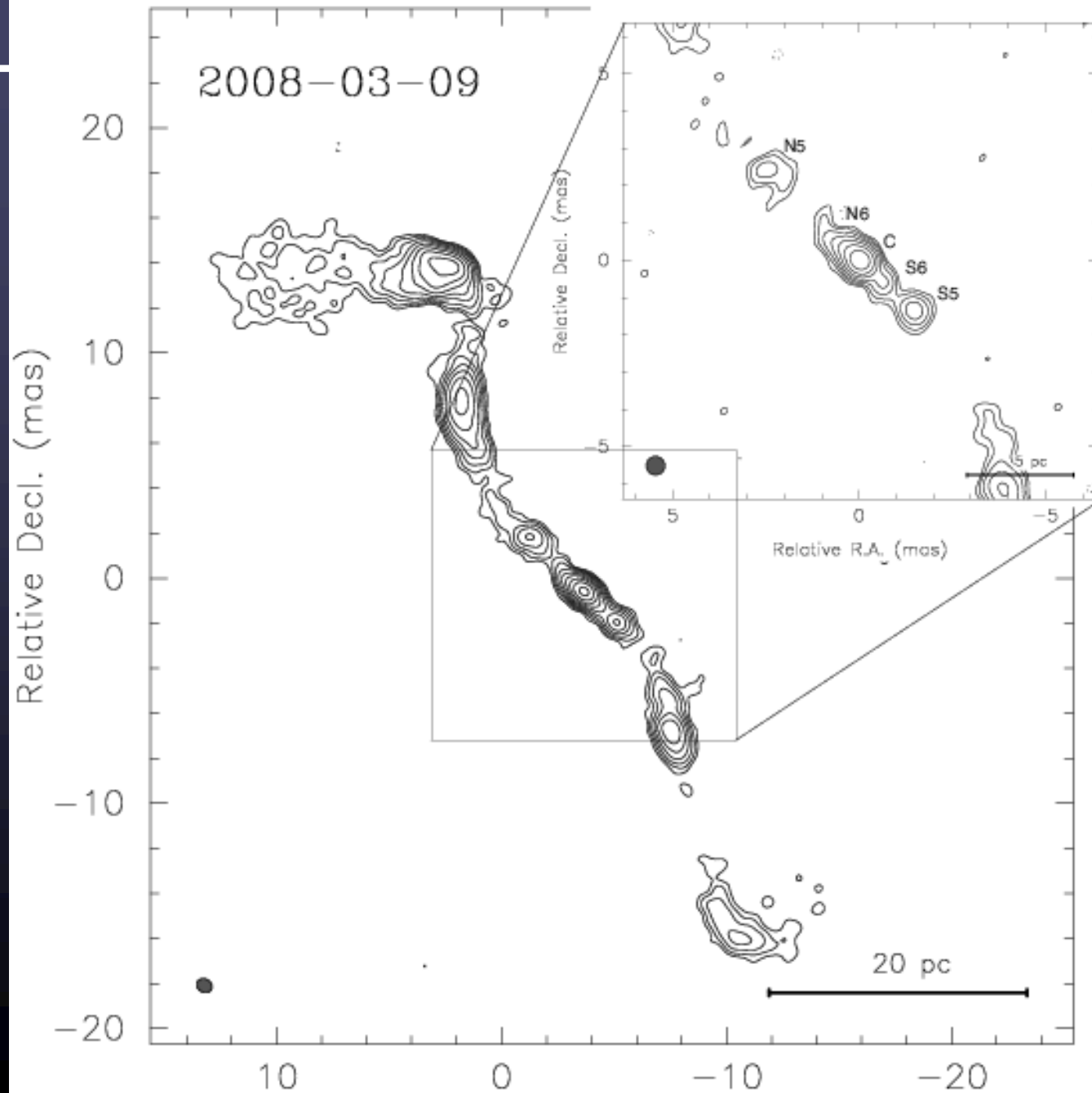
$$S_{\text{obs}} = S_{\text{exp}} e^{-\tau} \quad \tau = 3.2$$

$$EM = 1.7 \times 10^8 \text{ pc cm}^{-2} = n_e^2 \cdot 50 \text{ pc}$$

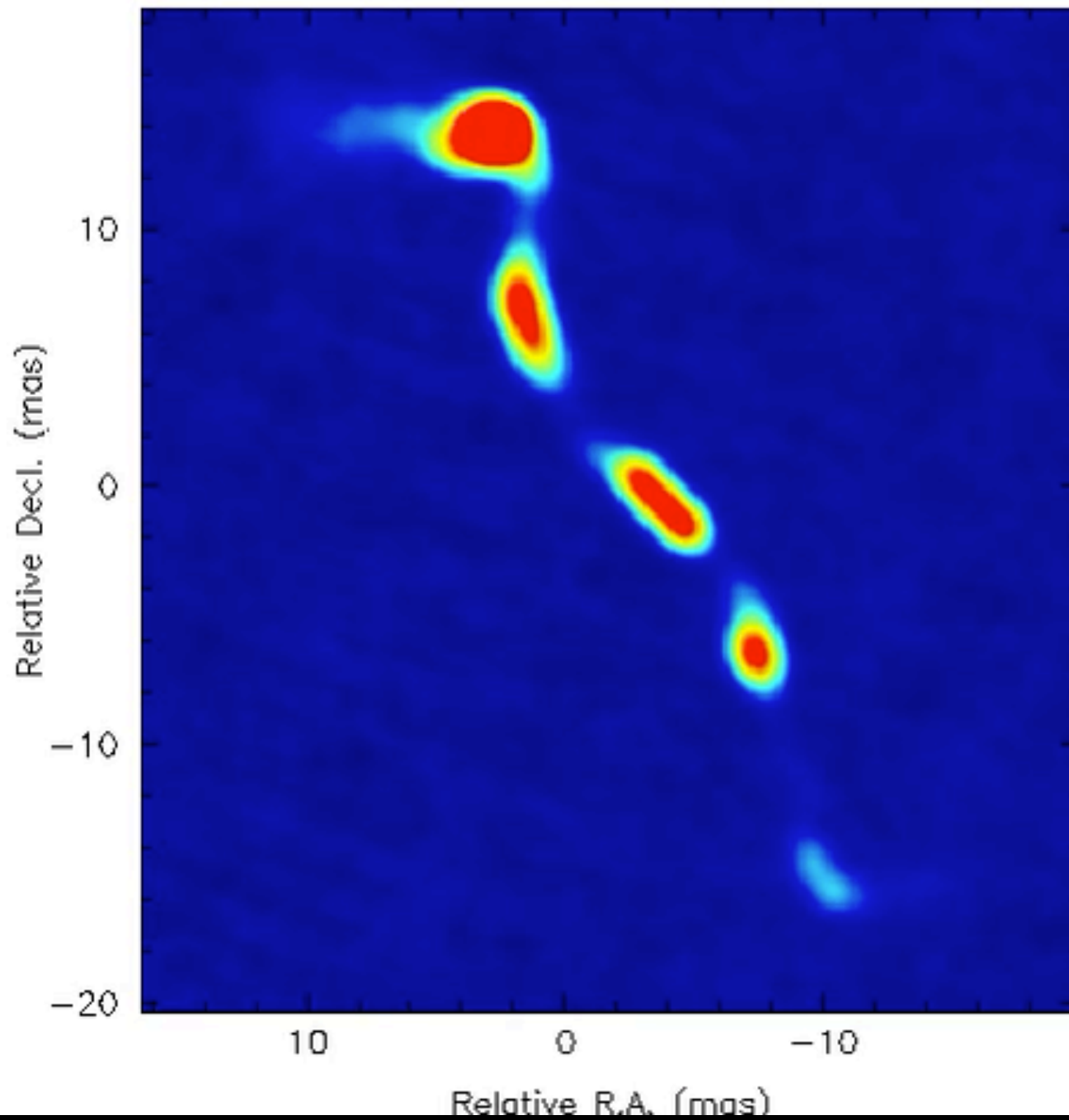


$$n_e = 1800 \text{ cm}^{-3}$$

Frequency: 8.421 GHz

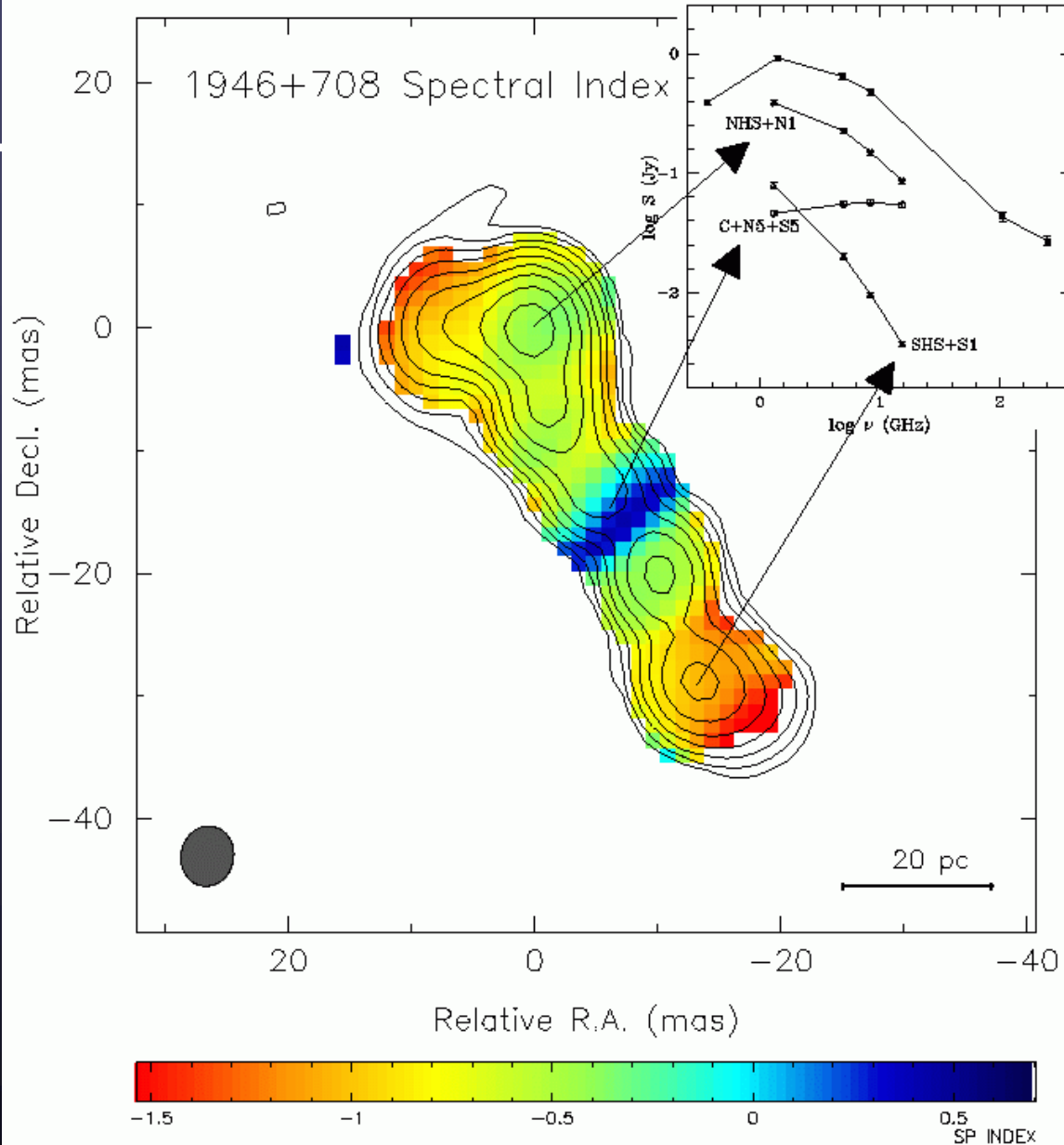


1946+708 VLBA 8.421 GHz 1995-03-22





# Free-free absorption in 1946+708



Peck & Taylor (2001)

Spectral index map from 1.3/5 GHz VLBI observations

free-free optical depth:

$$\tau_{\text{ff}} \sim T^{-3/2} n_e^2 \nu^{-2} d$$

$$N_e \sim 8 \times 10^{22} \text{ cm}^{-2}$$

ionization  $\sim 10\%$



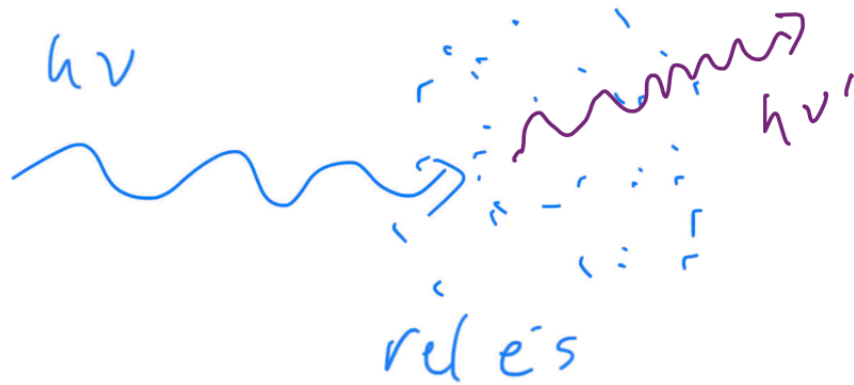
Free-Free Emission  
Polarization: Intrinsically zero

Beaming: Intrinsically none  
Radiation is isotropic



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## Bonus: Inverse Compton Emission



$$h\nu' = \gamma h\nu (1 + \beta \cos\theta)$$

$$-\frac{dE}{dt} = \sigma_T c u'_{\text{rad}}$$

$$u'_{\text{rad}} = \left[ \gamma \left( 1 + \frac{v}{c} \cos\theta \right) \right]^2 u_{\text{rad}}$$

$$-\frac{dE}{dt} = \frac{4}{3} \sigma_T c u_{\text{rad}} \gamma^2$$

looks a lot like synchrotron emission

$$-\frac{dE}{dt} = \frac{4}{3} \sigma_T c \gamma^2 u_{\text{mag}} \quad u_{\text{mag}} = B^2/8\pi$$

$$\frac{\frac{dE}{dt} \text{ Compton}}{\frac{dE}{dt} \text{ Syach}} = \frac{u_p h}{u_b}$$